

Chapters 22/23: Potential/Capacitance

Tuesday September 20th

Mini Exam 2 on Thursday:

Covers Chs. 21 and 22 (Gauss' law and potential)

Covers LONCAPA #3 to #6 (due this Wed.)

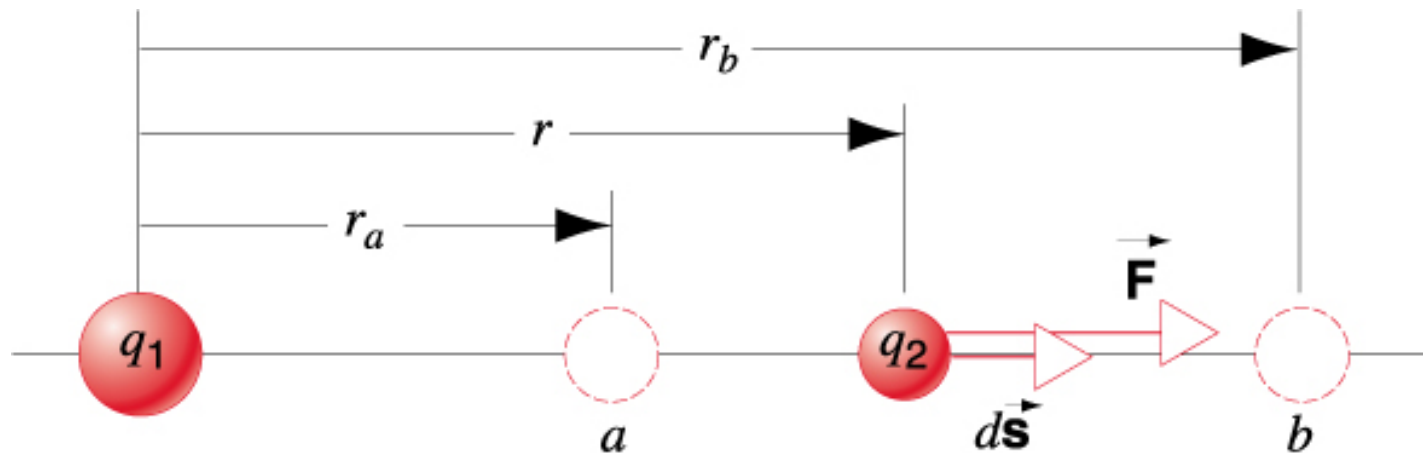
No formula sheet allowed!!

- Review: Electrostatic potential energy
- Review and continuation: Electrostatic potential
 - Relationship between V and E
- Capacitance
 - Definition
 - Examples
- Equipotential surfaces
 - Conductors
 - Relationship between E and V

Reading: up to page 386 in the text book (Chs. 22/23)

Electrostatic Potential Energy

- *The electrostatic (Coulomb) force is conservative.*
- *It is this property that allows us to define a scalar potential energy (one cannot do this for non conservative forces).*



$$\Delta U = - \int_a^b \vec{\mathbf{F}} \cdot d\vec{\mathbf{r}} = + \frac{1}{4\pi\epsilon_o} q_1 q_2 \left(\frac{1}{r_b} - \frac{1}{r_a} \right)$$

- *One can then apply energy conservation (PHY2048).*

Electrostatic Potential Energy

- *The potential energy is a property of both of the charges, not one or the other.*
- *If we choose a reference such that $U = 0$ when the charges are infinitely far apart, then we can simplify the expression for the potential energy as follows.*

$$U(r) = -\int_{\infty}^r \vec{\mathbf{F}} \cdot d\vec{\mathbf{r}} = +\frac{1}{4\pi\epsilon_0} q_1 q_2 \left(\frac{1}{r} - \frac{1}{\infty} \right) = \frac{1}{4\pi\epsilon_0} \frac{q_1 q_2}{r}$$

Again, the sign of U is not a problem. It is taken care of by the signs of the charges q_1 and q_2 .

The Electrostatic Potential

- We define a new quantity known as the *Electrostatic Potential* V , simply by dividing out the test charge q_o :

i.e.,

$$U = q_o V$$

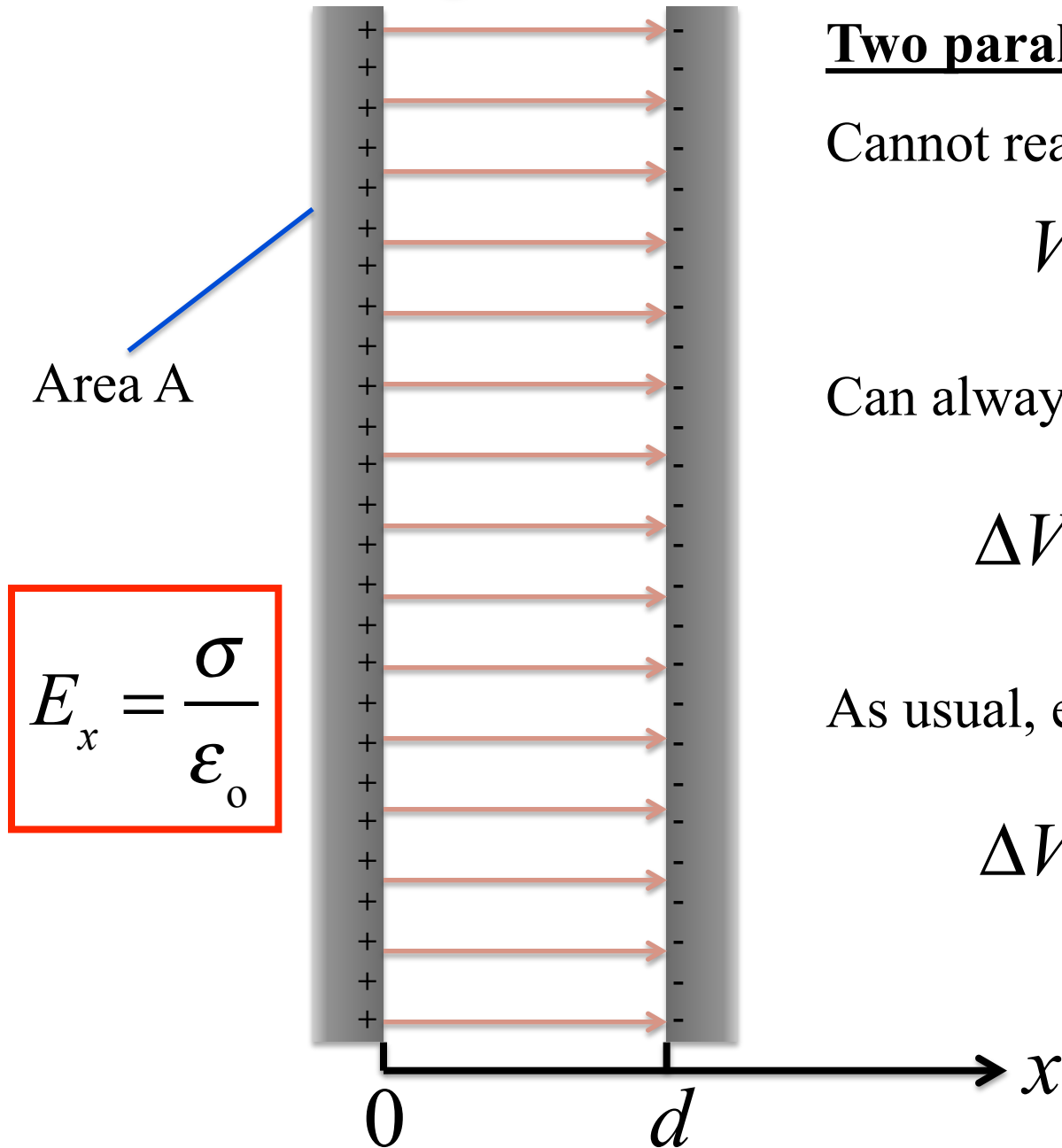
$$\Rightarrow \Delta V = - \int_a^b \vec{\mathbf{E}} \cdot d\vec{\mathbf{r}} = \frac{q}{4\pi\epsilon_o} \left(\frac{1}{r_b} - \frac{1}{r_a} \right)$$

- This ‘scalar potential’ depends only on the details of the source charge distribution (in this case, q).

$$V(r) = - \int_{\infty}^r \vec{\mathbf{E}} \cdot d\vec{\mathbf{r}} = \frac{q}{4\pi\epsilon_o} \frac{1}{r} = k \frac{q}{r}$$

- The absolute value of the potential is not important. As we shall see, it is only potential differences that really matter.

Calculating Potential Difference from E



$$E_x = \frac{\sigma}{\epsilon_0}$$

Two parallel conducting plates

Cannot really use:

$$V(r) = k \frac{q}{r}$$

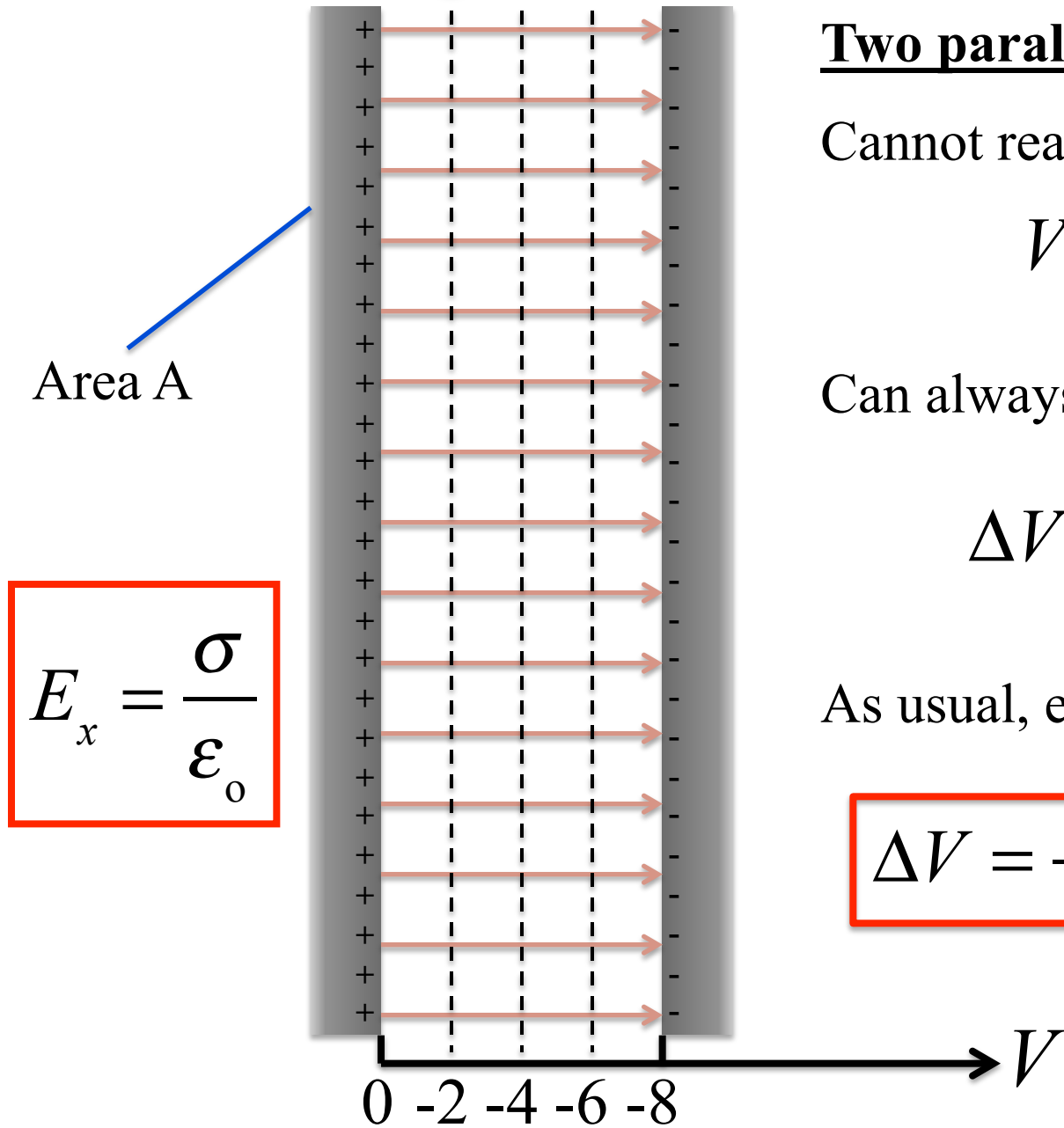
Can always use definition:

$$\Delta V = - \int_{r_1}^{r_2} \vec{E} \cdot d\vec{r}$$

As usual, exploit the symmetry:

$$\Delta V = - \int_0^x E_x dx$$

Calculating Potential Difference from E



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Two parallel conducting plates

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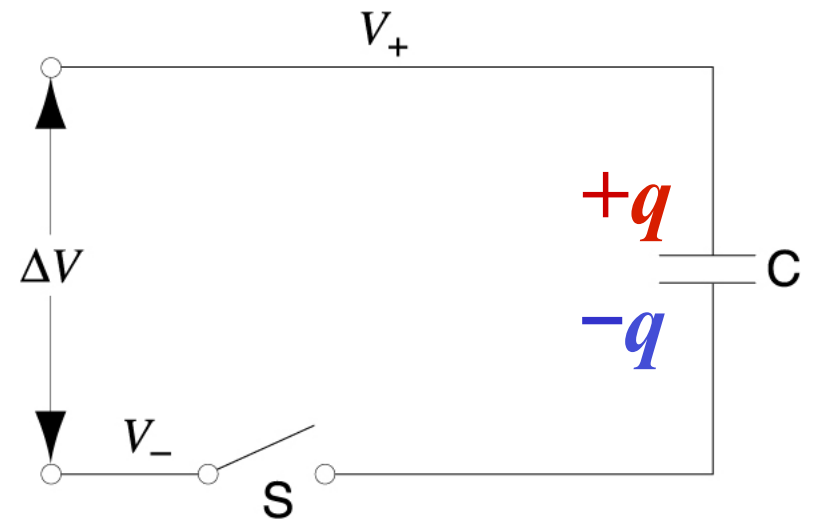
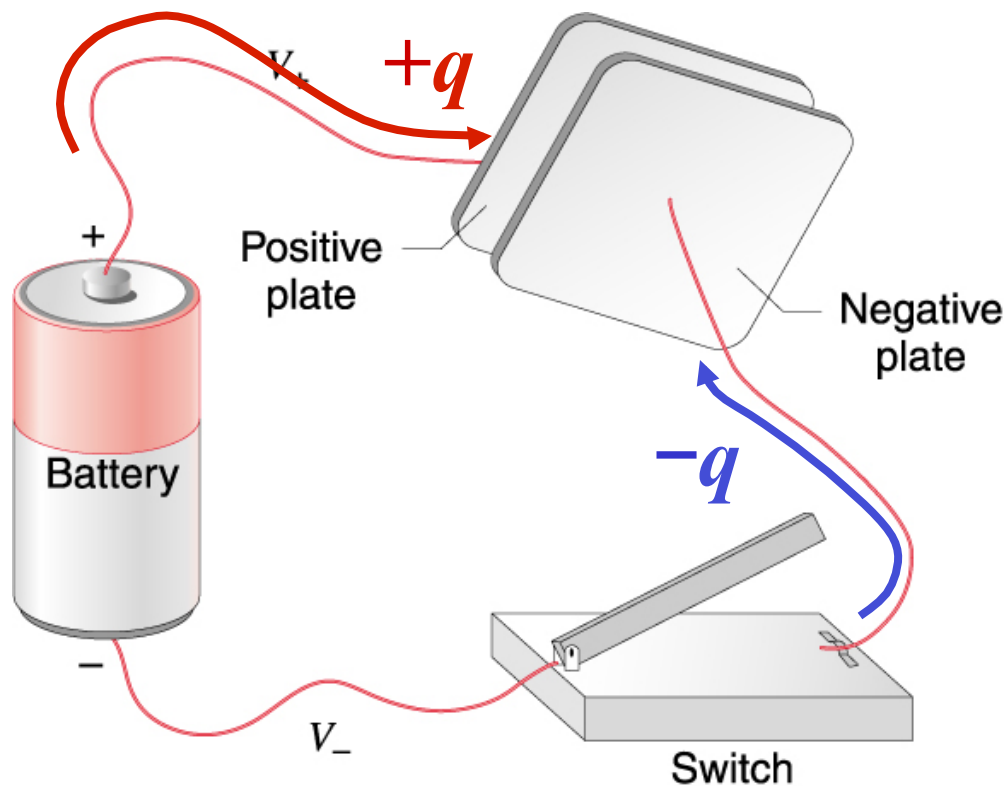
$$\Delta V = - \int_{r_1}^{r_2} \vec{E} \cdot d\vec{r}$$

As usual, exploit the symmetry:

$$\Delta V = -E_x \Delta x = -\frac{\sigma}{\epsilon_0} d$$

Capacitors

- Used to store energy in electromagnetic fields [in contrast to batteries (chemical cells) that store chemical energy].
- Capacitors can release electromagnetic energy much, much faster than chemical cells. They are thus very useful for applications requiring very rapid responses.



Capacitors

- The transfer of charge from one terminal of the capacitor to the other creates the electric field.
- Where there is a field, there must be a potential difference, i.e., a voltage difference between terminals.
- This leads to the definition of capacitance C :

$$Q = C\Delta V$$

- Q represents the magnitude of the excess charge on either plate. Another way of thinking of it is the charge that was transferred between the plates.

SI unit of capacitance: 1 farad (F) = 1 coulomb/volt

(after Michael Faraday)

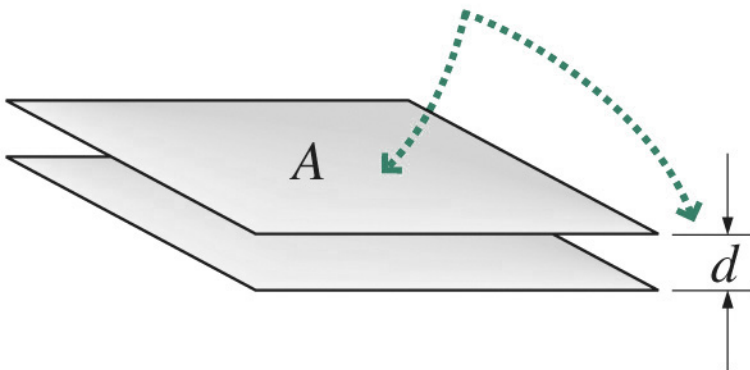
Capacitances more often have units of picofarad (pF) and microfarad (μ F)

Capacitors

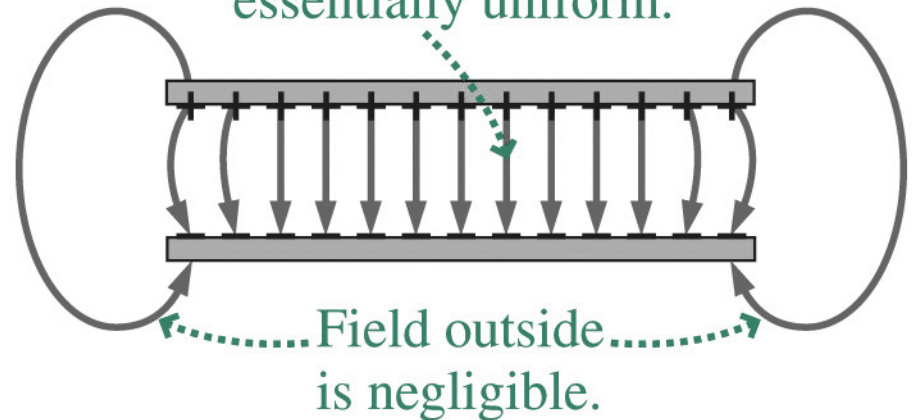
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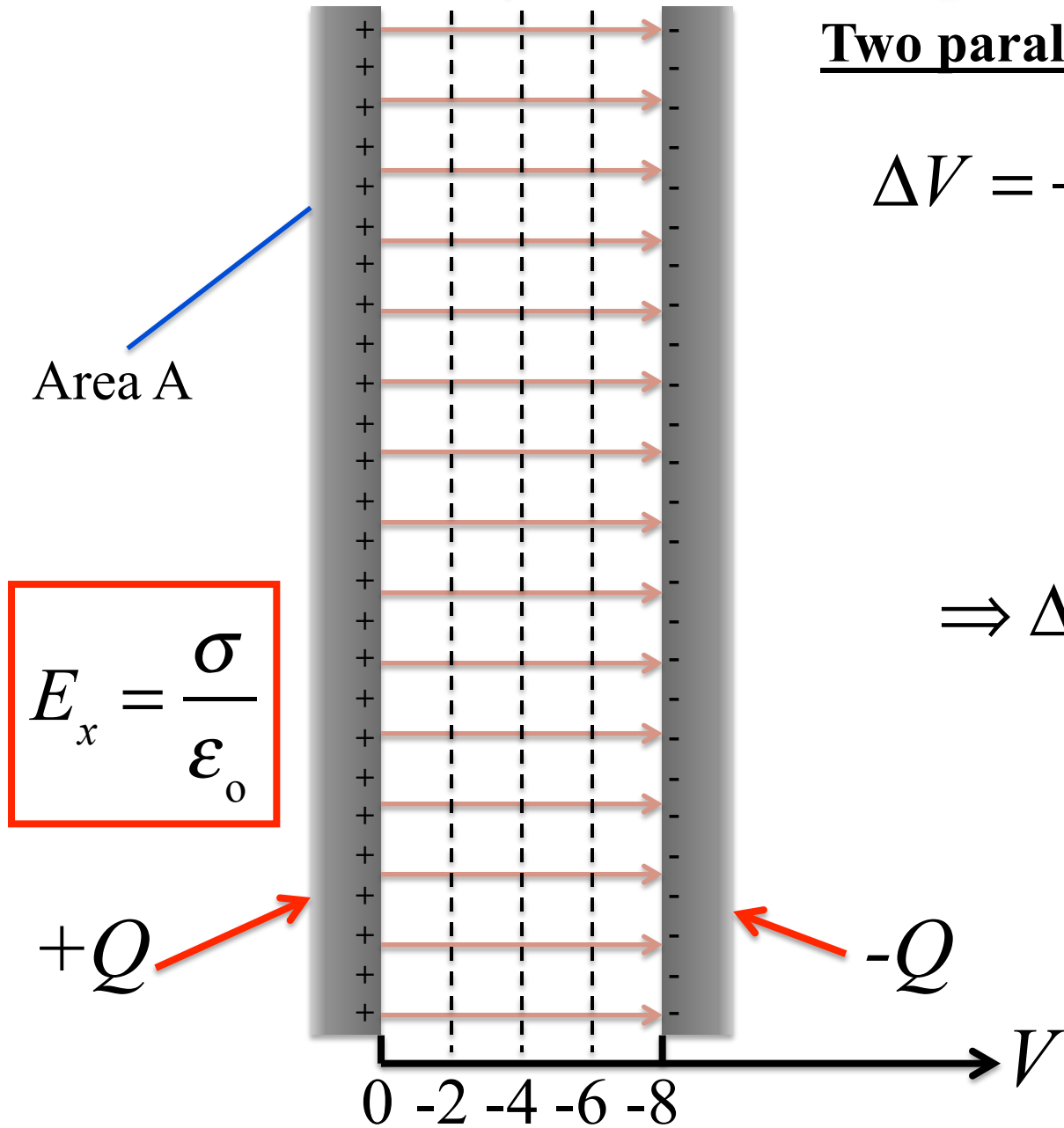
Conducting plates with area A are a small distance d apart.



Field inside is essentially uniform.



Our Example Involving Parallel Plates



$$E_x = \frac{\sigma}{\epsilon_0}$$

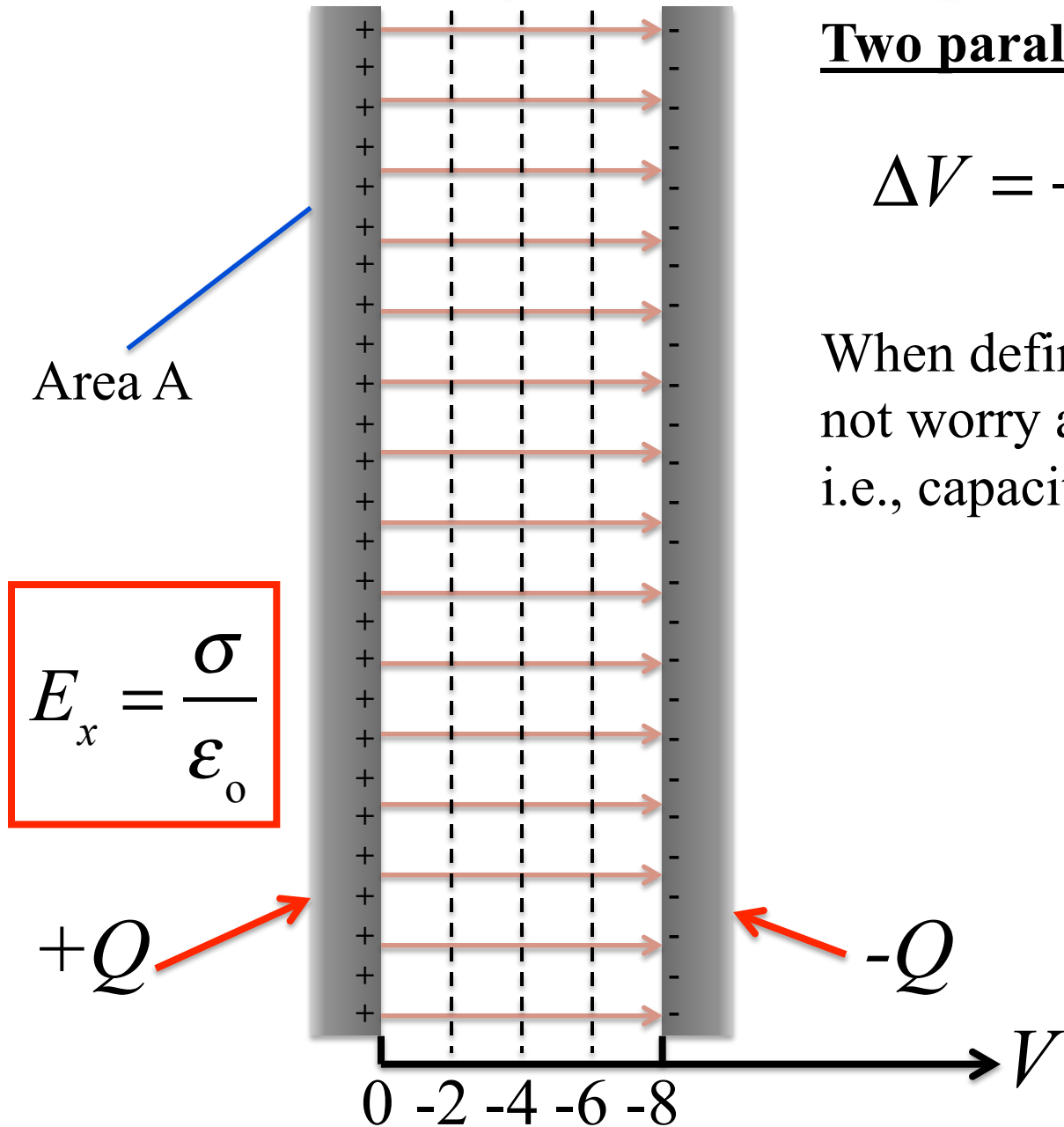
Two parallel conducting plates

$$\Delta V = -E_x \Delta x = -\frac{\sigma}{\epsilon_0} d$$

$$\sigma = \frac{Q}{A}$$

$$\Rightarrow \Delta V = -\frac{Q}{A\epsilon_0} d$$

Our Example Involving Parallel Plates



Two parallel conducting plates

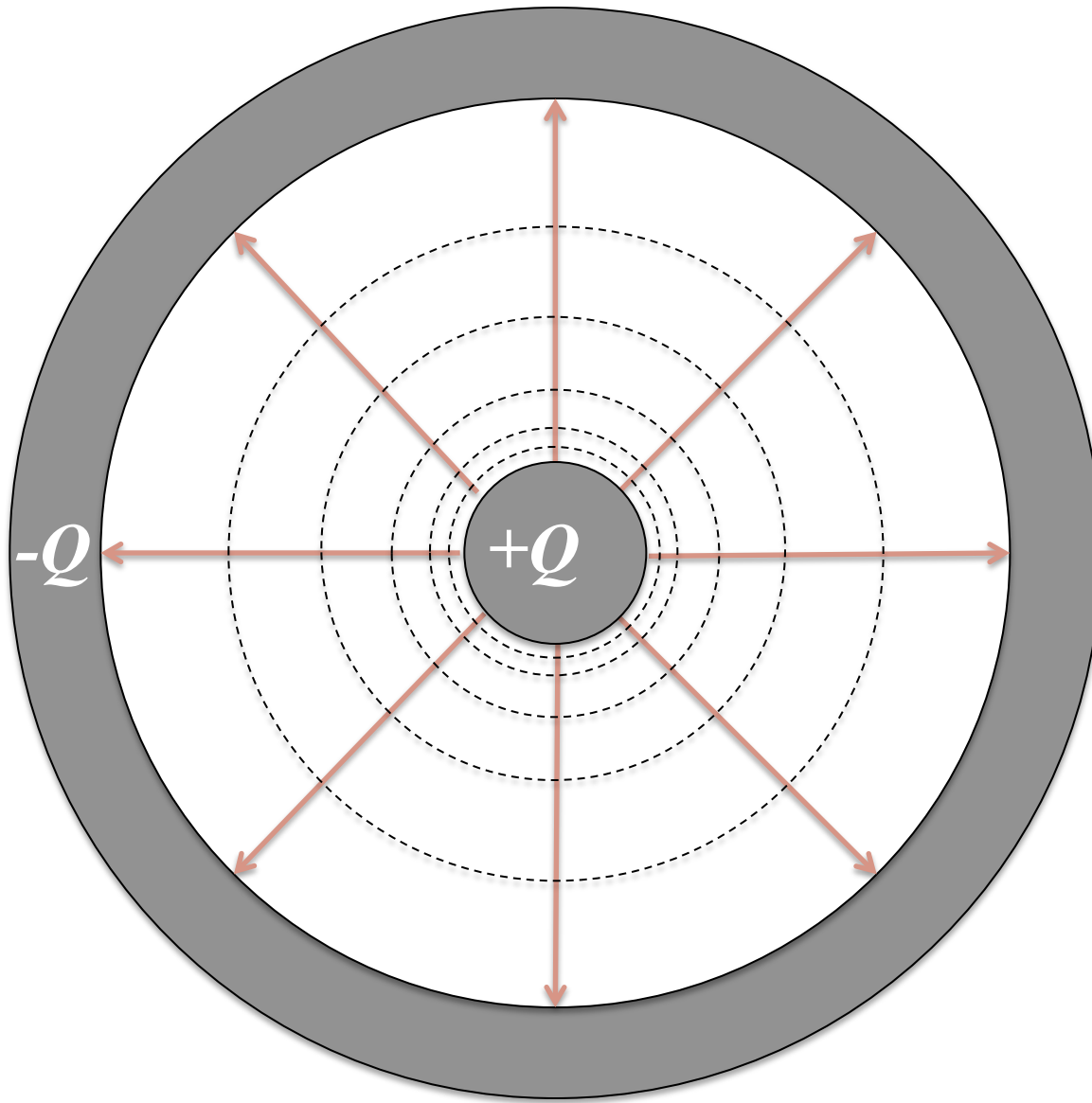
$$\Delta V = -E_x \Delta x = -\frac{\sigma}{\epsilon_0} d$$

When defining capacitance, we do not worry about sign of potential, i.e., capacitance is always positive

$$\Delta V = Q \frac{d}{A\epsilon_0} = \frac{Q}{C}$$

$$\Rightarrow C = \frac{\epsilon_0 A}{d}$$

Another Example: Concentric Spheres



Can use both

$$V(r) = k \frac{q}{r}$$

and

$$\Delta V = - \int_{r_1}^{r_2} \vec{\mathbf{E}} \cdot d\vec{\mathbf{r}}$$

$$\vec{\mathbf{E}}(r) = k \frac{q}{r^2} \hat{\mathbf{r}}$$

Some sophisticated vector calculus

The fundamental theorem of calculus:

$$\int_{x_1}^{x_2} \left(\frac{df}{dx} \right) dx = f(x_2) - f(x_1)$$

Recall the 1D example involving the parallel plates:

$$\Delta V = - \int_{r_1}^{r_2} \vec{\mathbf{E}} \cdot d\vec{\mathbf{r}} = \int_{x_1}^{x_2} (-E_x) dx = V_2 - V_1$$

$$\Rightarrow E_x = - \frac{dV}{dx} \approx - \frac{\Delta V}{\Delta x}$$

Potential Does Not Kill You



Some more sophisticated vector calculus

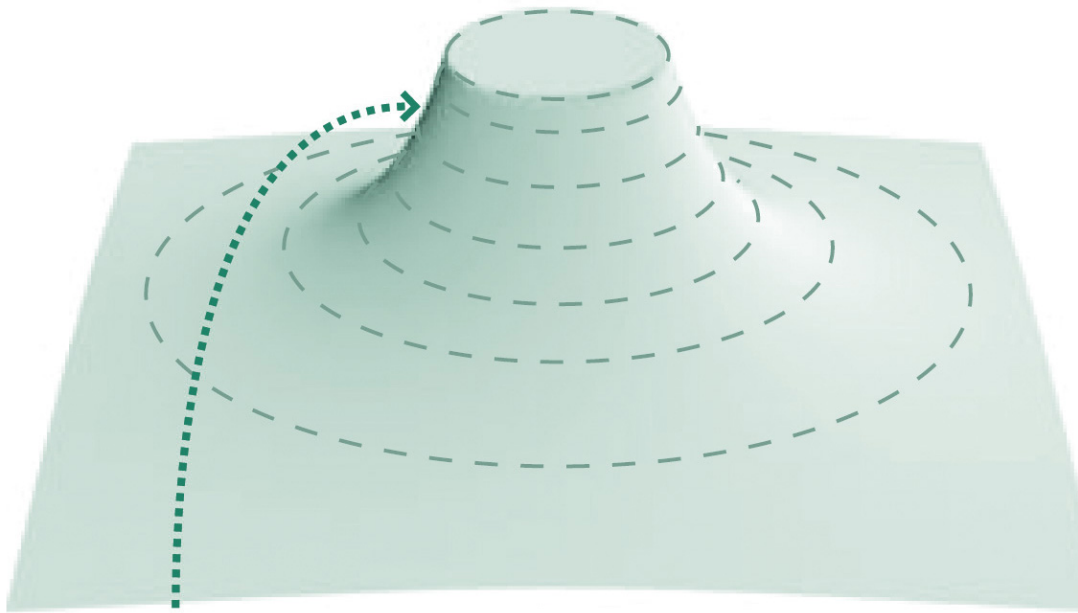
The fundamental theorem of calculus:

In 3D:

$$\Delta V = \int_{\vec{r}_1}^{\vec{r}_2} \nabla V \cdot d\vec{r} = V(\vec{r}_2) - V(\vec{r}_1)$$

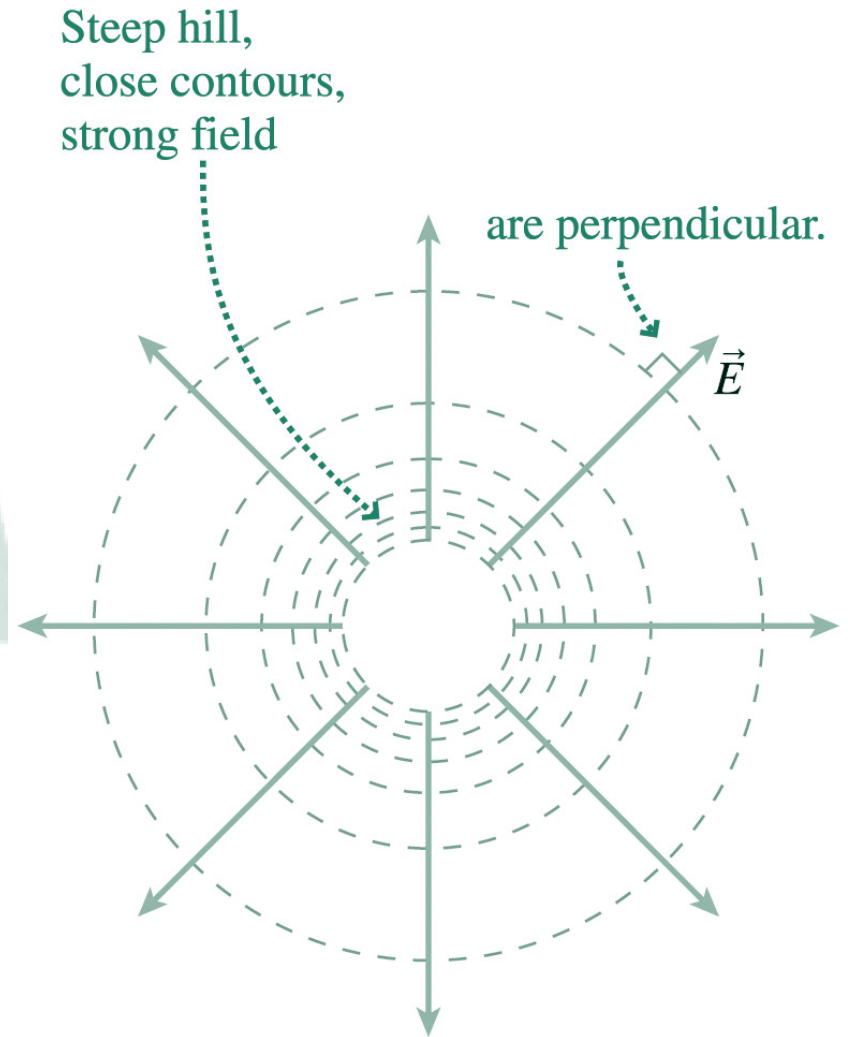
$$\vec{E} = -\nabla V = -\left(\frac{\partial V}{\partial x} \hat{i} + \frac{\partial V}{\partial y} \hat{j} + \frac{\partial V}{\partial z} \hat{k} \right)$$

Equipotential lines/surfaces



Steep hill,
close contours,
strong field

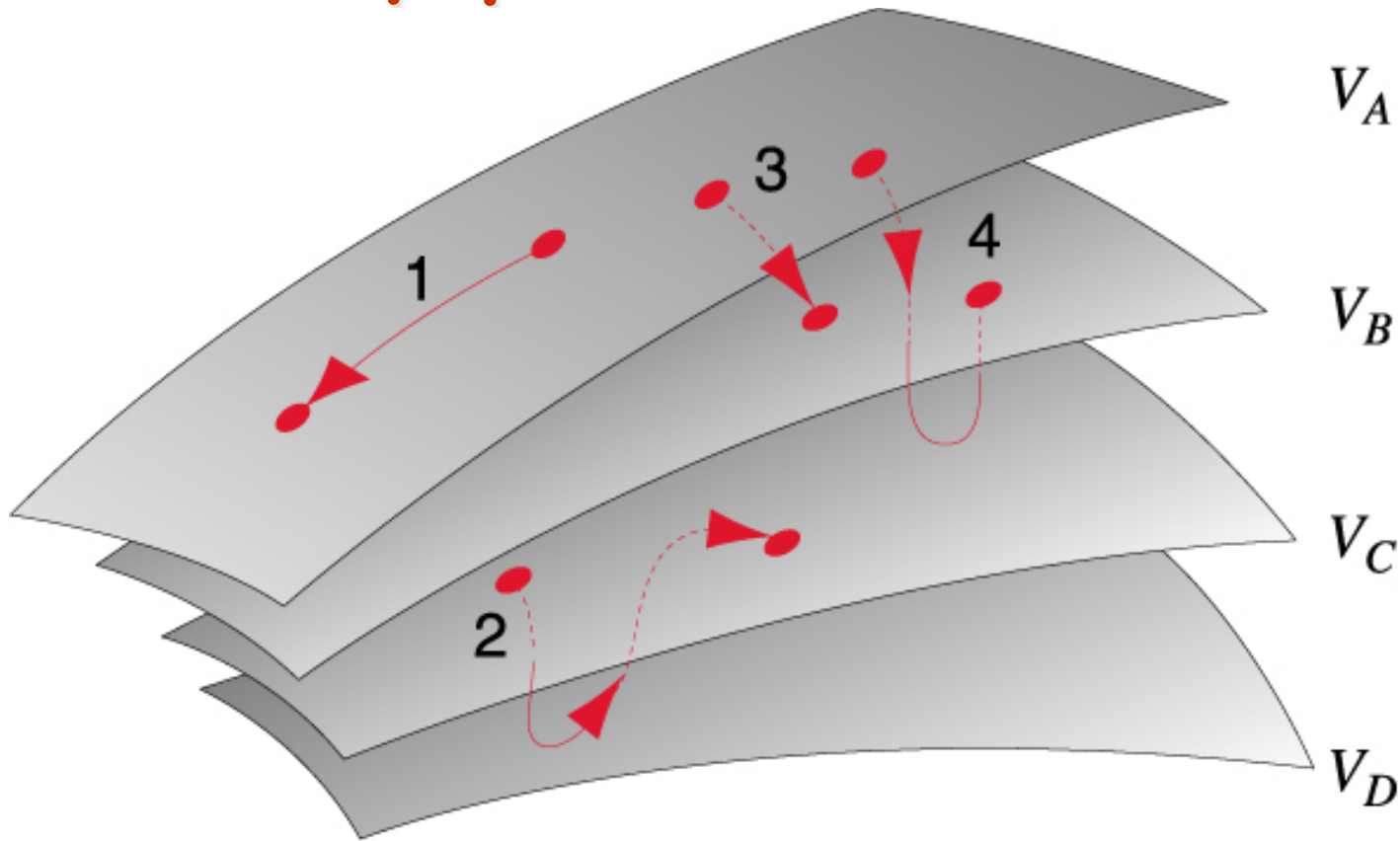
$$E_r \approx -\frac{\Delta V}{\Delta r}$$



Steep hill,
close contours,
strong field

are perpendicular.

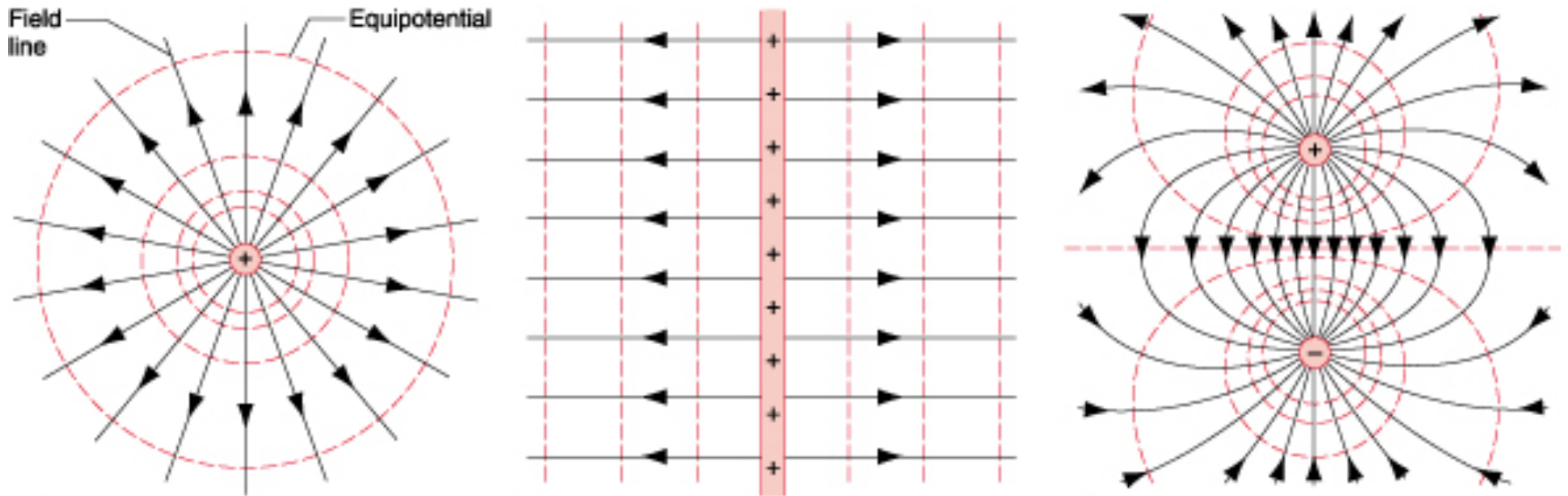
Equipotential surfaces



$$W_{field} = -q\Delta V = -q \int_a^b \vec{\mathbf{E}} \cdot d\vec{\mathbf{r}} = -q(V_b - V_a)$$

- ***E-field lines perpendicular to equipotential surfaces***
- ***Field does positive work when q accelerated by field***

Equipotential surfaces



Here are some example surfaces including field lines (point charge, infinite charged plane and a dipole).

- By spacing the equipotential surfaces by the same potential difference (ΔV), one can get a feel for the electric field strength ($E = -dV/dr$), i.e. the closer the spacing, the stronger the field.*